

Lecture VII

Hilbert space fibre bundle
tenfold way + examples

previous lecture T, P, C symmetries & properties

T : time reversal P : particle-hole C : chiral symmetry
if $T^2 = -1 \rightarrow$ Kramers deg. ϵ \rightarrow twice degenerate ϵ \rightarrow sym. spectrum $\mathcal{K} = \begin{pmatrix} 0 & \mathbb{Q} \\ \mathbb{Q}^+ & 0 \end{pmatrix}$

Hilbert subspace fibre bundle cont.
repet. $E \subseteq \mathcal{H} \otimes B$ is the total space of fibre bundle (inherits topology) as subspace of product $E \xrightarrow{\pi} B$ with $\pi: E \rightarrow B$ $(\varphi, \nu) \mapsto \varphi$
 $\mathcal{K}(\varphi, \nu)$ defined by Hamiltonian $\mathcal{K}(\varphi)$ as eigenvalues & eigenvectors are continuous.
 $\mathcal{K}(\varphi) \in \mathbb{C}^n \times \mathbb{R}^n$ locally parameter manifold $\varphi \in B$
 $\varphi \in B$ there is a neighborhood $U \subseteq B$ with $U \cong \mathbb{R}^n$

π surjective: $\pi(E) = \{ \pi(\varphi, \nu) \mid \varphi \in B \} = \{ \varphi \mid \varphi \in B \} = B$

π continuous: By definition of the product space topology of $\mathcal{H} \otimes B$.

local triviality: For $\varphi \in B$, take neighborhood $U \subseteq B$ with $U \cong \mathbb{R}^n$. Observe that $F(\varphi')$ with $\varphi' \in U$ is homeomorphic to \mathbb{C}^n by $\tilde{\nu} \in F(\varphi) \mapsto \tilde{\nu} \cdot \sum_{j=1}^n \tilde{\nu}_j(\varphi) |j\rangle \in \mathbb{C}^n$. $F(\varphi') = \pi^{-1}(\varphi')$

Then $i: U \times \mathbb{C}^n \rightarrow \pi^{-1}(U)$ with $(\varphi, \tilde{\nu}) \mapsto (\varphi, \sum_{j=1}^n \tilde{\nu}_j(\varphi) |j\rangle)$ is a homeomorphism with $\pi(i(U \times \mathbb{C}^n)) = \pi(\{ \varphi' \mid \sum_{j=1}^n \tilde{\nu}_j(\varphi) |j\rangle \mid \varphi' \in U \}) = \{ \varphi' \mid \varphi' \in U \} = U$

Question: What happens if $\mathcal{K}(\varphi)$ discontinuous

Topological classification

Topological classification differs for cont theories & crystal (lattice) theories | Ryu, Schnyder, Furusaki

$\mathcal{H} = \int_{\mathbb{R}^d} dx \psi^\dagger(x) \mathcal{H} \psi(x) + \text{anomalous terms}$
 $= \int_{\mathbb{R}^d} dk \tilde{\psi}^\dagger(k) \tilde{\mathcal{H}}(k) \tilde{\psi}(k) + \text{anomalous terms}$
 $\mathbb{S}^d = \mathbb{R}^d / \infty$ [$\tilde{\mathcal{H}}(k) \rightarrow 0$ if $k \rightarrow \infty$ (low-energy theory)]
 \rightarrow identify all points at $k \rightarrow \infty$

first focus on the cont. low-energy theories

Important equivalence classes for top. classification
- orders of occ & unocc bands do not matter
 \rightarrow equivalence classes $U(n+m)/\sim (U(n) \times U(m)) \cong \mathbb{Z}$

Ultimately $Q: \mathbb{S}^d \rightarrow G(n+m, m)$
Topological classification of Hamiltonian by $\Pi_d(G(n+m, m))$

Cartan	P	T	C	1	2	3	4
A	0	0	0	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}
AIII	0	0	1	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}
AI	1	0	0	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}
BOI	1	1	1	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}
D	-1	1	1	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}
DIII	0	1	0	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}
AII	-1	0	0	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}
CII	-1	-1	1	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}
C	0	-1	0	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}
CI	1	-1	1	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}

\circ : (low-energy) strong topological invariants.
 \square : weak topological invariants.
Remark: Additional P, T, C symmetries change the target manifold. [We do not start with all unitary top. Hamiltonians] \rightarrow different $\Pi_d \Rightarrow$ different top. classification

Example: DIII material, top. index \mathbb{Z} , weak 1D & 2D top indices

Examples class A, 2D $\Rightarrow \mathbb{Z}$ invariant Quantum Hall effect.
 $\mathbb{Z} = \frac{1}{2\pi i} \int d^2k \text{Tr} \{ P_k [d_1 P_k d_2 P_k d_3 P_k] P_k \}$
TKNM invariant 1362
 $\mathbb{Z} \hat{=} \text{number of edge modes.}$

Kitaev chain 1D top. super conductor in D with \mathbb{Z}_2 invariant

Example next time.