

# Topology in condensed matter physics

## Exercise sheet 5

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### 5.1 Bogoliubov transformations

A Bogoliubov transformation is used to diagonalize a second quantized, quadratic Hamiltonian with anomalous terms. These are terms of the form creation operator times creation operator and annihilation operator times annihilation operator, which appear in the mean-field description of superconductors and in the solution of the interacting fermion problem in one dimension.

Given the annihilation operators  $c_i$  for  $i$  in a countable index set  $I$ , we introduce the transformed annihilation operator

$$d_j = \sum_{i \in I} A_{j,i} c_i + B_{j,i} c_i^\dagger, \quad (1)$$

for each  $j \in I$ .

1. Assume that the operators  $c_i$  for  $i \in I$  are fermionic, i.e.,  $\{c_i, c_j\} = 0$ ,  $\{c_i^\dagger, c_j\} = \delta_{i,j}$ . Derive (necessary and sufficient) conditions for the matrices  $A$  and  $B$  such that the  $d_j$  are fermionic as well.
2. Assume that the operators  $c_i$  for  $i \in I$  are bosonic, i.e.,  $[c_i, c_j] = 0$ ,  $[c_i, c_j^\dagger] = \delta_{i,j}$ . Derive (necessary and sufficient) conditions for the matrices  $A$  and  $B$  such that the  $d_j$  are bosonic as well.
3. Find a parametrization for the matrices  $A$  and  $B$  of the Bogoliubov transformation of a single bosonic mode (i.e.,  $I = \{1\}$ ).
4. (\* points) Diagonalize

$$H = \alpha b^\dagger b + b^\dagger b^\dagger + bb, \quad (2)$$

where  $\alpha \in \mathbb{R}$  and  $b$  is a bosonic annihilation operator, by finding a boson  $d$  such that

$$H = \epsilon d^\dagger d + \text{const.} \quad (3)$$

For which values of  $\alpha$  is such a diagonalization impossible?

*Hint: Use  $[d, H] = \epsilon d = \epsilon(Ab + Bb^\dagger)$  to obtain linear equations for the coefficients  $A$  and  $B$ .* Interpretation: If the Hamiltonian is not diagonalizable by a Bogoliubov transformation, it describes an unstable equilibrium.

### 5.2 Symmetry reduction of Hamiltonians

Consider the following set of Hamiltonians, a two-site tight-binding model for spinful fermions, for instance, describing the valence electrons of a binary molecule in a magnetic field

$$\mathcal{H} = BS_x + t \sum_{\sigma \in \{\uparrow, \downarrow\}} c_{1,\sigma}^\dagger c_{2,\sigma} + c_{2,\sigma}^\dagger c_{1,\sigma}, \quad (4)$$

with the real parameters  $t, B \in \mathbb{R}$  describing electron transfer between the two sites and the magnetic field in  $x$ -direction, respectively, and  $S_x = \frac{\hbar}{2} \sum_{i=1}^2 c_{i,\uparrow}^\dagger c_{i,\downarrow} + c_{i,\downarrow}^\dagger c_{i,\uparrow}$  being the total  $x$ -spin of the particles.

1. Show that  $\mathcal{H}$  conserves the total  $x$ -spin, i.e.,  $[\mathcal{H}, S_x] = 0$ . Does that remain true for an inhomogeneous magnetic field in  $x$ -direction?
2. What is the corresponding unitary symmetry of the family of Hamiltonians?
3. Derive the single particle Hamiltonian by writing  $\mathcal{H}$  in the form

$$\sum_{a,b=1}^4 C_a^\dagger h_{a,b} C_b. \quad (5)$$

with  $c = (c_{1,\uparrow}, c_{2,\uparrow}, c_{1,\downarrow}, c_{2,\downarrow})$ . Here,  $h$  is the single particle Hamiltonian.

4. Block diagonalizes the Hamiltonian by rotating the spin basis into  $x$ -direction. Can you represent this transformation and the result in the matrix formulation introduced in Eq. (5)? *Hint: This can be done by introducing the rotated quasi-particles  $\tilde{c}_{i,\uparrow} = \frac{1}{\sqrt{2}}(c_{i,\uparrow} + c_{i,\downarrow})$  and  $\tilde{c}_{i,\downarrow} = \frac{1}{\sqrt{2}}(-c_{i,\uparrow} + c_{i,\downarrow})$ . (This has an insightful explanation. Which?)*

**End**